Note on Coupled Linear Systems Related to Two Soliton Collision for the Quartic gKdV Equation

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ABSTRACT

We consider linear systems related to the description of the collision of two solitons for the gKdV equation with quartic nonlinearity. The computations presented in this note are applied in Martel and Merle [10] to prove a result concerning inelastic (but almost elastic) collision for a nonintegrable equation.

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Introduction

We consider the quartic generalized Korteweg-de Vries (gKdV) equation:

$$\partial_t u + \partial_x (\partial_x^2 u + u^4) = 0, \quad x, t \in \mathbb{R}.$$
 (1)

The objective of this note is to solve some linear coupled systems related to the collision of two solitons for (1) (see also section 2 for the case where the nonlinearity is u^2). Recall that equation (1), unlike the corresponding equations with nonlinearities u^2 or u^3 , is not integrable. In particular, the collision of two solitary waves of (1)

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is not described by explicit 2-soliton solutions. In [10], we have introduced a new method to understand the collision of two solitary waves of (1) in the special case where one solitary wave is small with respect to the other. More precisely, denote by Q the unique even solution of

$$Q > 0$$
, $Q'' + Q^4 = Q$, $Q \in H^1(\mathbb{R})$, i.e., $Q(x) = \left(\frac{5}{2\cosh^2(\frac{3}{2}x)}\right)^{\frac{1}{3}}$, (2)

and, for any c > 0, let $Q_c(x) = c^{\frac{1}{3}}Q(\sqrt{c}x)$ be a solution of $Q''_c + Q^4_c = cQ_c$. Then, for any $x_0 \in \mathbb{R}$, c > 0, the functions

$$R_{c,x_0}(t,x) = Q_c(x - x_0 - ct)$$

are solitary wave solutions of (1), called *solitons*.

From [7] (see also [8]), there exist solutions u(t,x) of (1) which are asymptotic N-soliton solutions at $t \to -\infty$ in the following sense: let $N \ge 1$, $c_1 > \cdots > c_N > 0$, and $x_1, \ldots, x_N \in \mathbb{R}$, there exists a unique H^1 solution U of (1) such that

$$\lim_{t \to -\infty} \left\| U(t) - \sum_{j=1}^{N} Q_{c_j}(. - x_j - c_j t) \right\|_{H^1(\mathbb{R})} = 0.$$

The behavior displayed by these solutions is stable in some sense. Considering for example the case of two solitons, there exist a large class of solutions such that, as $t \sim -\infty$,

$$u(t,x) = Q_{c_1}(x-x_1-c_1t) + Q_{c_2}(x-x_2-c_2t) + \eta(t,x),$$

where $c_1 > c_2$ and $\eta(t)$ is a dispersion term small in the energy space H^1 with respect to Q_{c_1} , Q_{c_2} (see [8]). From the Physics point of view, the two solitons Q_{c_1} and Q_{c_2} have to collide at some time t_0 . In the special case $c_2 \ll c_1$ (or equivalently, $\|Q_{c_2}\|_{H^1} \ll \|Q_{c_1}\|_{H^1}$) and $\|\eta(t)\|_{H^1} \ll \|Q_{c_2}\|_{H^1}$, for t close to $-\infty$, we have introduced in [10] explicit computations allowing to understand the collision at the main orders, using a new nonlinear "basis" to write and compute an approximate solution v(t,x) up to any order of size.

Recall that the problem of collision of two solitons is a classical question in nonlinear wave propagation (see [3,13,14]). In the so-called integrable cases (i.e., when the nonlinearity in (1) is u^2 or u^3) it is well-known that there exist explicit multi-soliton solutions, describing the elastic collision of several solitons (see Hirota [4], Lax [5], Wadati and Toda [12], and the review paper Miura [11]). Note that in experiments, or numerically for more accurate nonintegrable models (see Craig et al. [2], Li and Sattinger [6], Bona et al. [1], and other references in [10]), this remarkable property is mainly preserved, i.e., the collision of two solitons is almost elastic; however, a (very small) residual part is observed after the collision. Equation (1) being not integrable, explicit N-soliton solutions are not available in this case. The results obtained in [10],

using Theorem A of the present paper, are the first rigorous results concerning inelastic (but almost elastic) collision in a nonintegrable situation. We refer to the introduction and the references in [10] for a overview on these questions.

The approximate solution v(t,x) in [10] has the following structure: let $k_0 \ge 1$, $\ell_0 \ge 0$, and

$$\Sigma_0 = \{(k, \ell), 1 \le k \le k_0, 0 \le \ell \le \ell_0\}.$$

We set

$$y_c = x + (1 - c)t$$
 and $R_c(t, x) = Q_c(y_c),$
 $y = x - \alpha(y_c)$ and $R(t, x) = Q(y),$

where for $(a_{k,\ell})_{(k,\ell)\in\Sigma_0}$,

$$\alpha(s) = \int_0^s \beta(s')ds', \qquad \beta(s) = \sum_{(k,\ell) \in \Sigma_0} a_{k,\ell} c^{\ell} Q_c^k(s).$$

The form of v(t,x) is

$$v(t,x) = Q(y) + Q_c(y_c) + W(t,x),$$

$$W(t,x) = \sum_{(k,\ell) \in \Sigma_0} c^{\ell} (Q_c^k(y_c) A_{k,\ell}(y) + (Q_c^k)'(y_c) B_{k,\ell}(y)),$$

where $a_{k,\ell}$, $A_{k,\ell}$, $B_{k,\ell}$ are to be determined. Let $S(t) = \partial_t v + \partial_x (\partial_x^2 v - v + v^4)$. We have obtained in [10, Proposition 2.1]:

$$S(t,x) = -\sum_{(k,\ell)\in\Sigma_0} c^{\ell} Q_c^k(y_c) \Big((\mathcal{L}A_{k,\ell})' + a_{k,\ell} (3Q - 2Q^4)' - F_{k,\ell} \Big) r(y)$$

$$- \sum_{(k,\ell)\in\Sigma_0} c^{\ell} (Q_c^k)'(y_c) \Big((\mathcal{L}B_{k,\ell})' + a_{k,\ell} (3Q'') - (3A''_{k,\ell} + 4Q^3 A_{k,\ell}) - G_{k,\ell} \Big) (y)$$

$$+ \mathcal{E}(t,x),$$

where

$$\mathcal{L}w = -\partial_x^2 w + w - 4Q^3 w,$$

 $\mathcal{E}(t,x)$ is a lower order term and the functions $F_{k,\ell}$, $G_{k,\ell}$ are second member terms in the sense that they depend only on $(a_{k',\ell'},A_{k',\ell'},B_{k',\ell'})$, for (k',ℓ') such that k' < k, $\ell' \le \ell$, or $k' \le k$, $\ell' < \ell$. Therefore, to obtain a suitable approximate solution (i.e., for which $S = \mathcal{E}$), we want to solve by induction on (k,ℓ) the following systems:

$$(\Omega_{k,\ell}) \begin{cases} (\mathcal{L}A_{k,\ell})' + a_{k,\ell}(3Q - 2Q^4)' = F_{k,\ell} \\ (\mathcal{L}B_{k,\ell})' + a_{k,\ell}(3Q'') - 3A_{k,\ell}'' - 4Q^3A_{k,\ell} = G_{k,\ell}. \end{cases}$$

We have proved in [10] that for a certain structure of $F_{k,\ell}$, $G_{k,\ell}$, one could always solve the system (Ω) , see [10, Propositions 2.2 and 2.3].

Using this approximate solution v, and analysis arguments (allowing to relate v(t) to solutions u(t) of (1)), we have proved the following result in [10, Theorem 1].

Non existence of a pure 2-soliton [10]. Let c > 0. There exists $\epsilon_0 > 0$ such that, if $c < \epsilon_0$, then there exists no pure 2-soliton solution with speeds 1 and c. More precisely, let u(t) be the unique H^1 solution of (1) such that

$$\lim_{t \to -\infty} ||u(t) - Q(\cdot - t) - Q_c(\cdot - ct)||_{H^1} = 0.$$

Then, there exist $x_1^+, x_2^+, c_1^+ > c_2^+ > 0$, and K > 0 such that

$$w^{+}(t,x) = u(t,x) - Q_{c_{1}^{+}}(x - x_{1}^{+} - c_{1}^{+}t) - Q_{c_{2}^{+}}(x - x_{2}^{+} - c_{2}^{+}t)$$

satisfies

$$\lim_{t \to +\infty} \|w^{+}(t)\|_{H^{1}(x > \frac{1}{10}ct)} = 0,$$

$$\frac{1}{K}c^{\frac{17}{6}} \le c_{1}^{+} - 1 \le Kc^{\frac{11}{12}}, \qquad \frac{1}{K}c^{\frac{8}{3}} \le 1 - \frac{c_{2}^{+}}{c} \le Kc^{\frac{1}{3}}, \tag{3}$$

$$\frac{1}{K}c^{\frac{17}{12}} < \|\partial_{x}w^{+}(t)\|_{L^{2}} + \sqrt{c_{1}c} \|w^{+}(t)\|_{L^{2}} < Kc^{\frac{11}{12}}, \qquad for \ t \ large. \tag{4}$$

The lower bound in (4) gives the minimal distance of u(t) to a two soliton solution after the collision. It is thus a qualitative version of the nonexistence of pure 2-soliton solution. Estimate (3) means that the speed of faster soliton is increased and the speed of the smaller one is decreased due to the collision.

For the proof of this result in [10], we needed in particular the resolution of $(\Omega_{1,0})$ and the value of $b_{2,0} = \lim_{x \to +\infty} B_{2,0}(x)$ in $(\Omega_{2,0})$. Let \mathcal{Y} be the set of functions $f \in C^{\infty}(\mathbb{R})$ such that

$$\forall j \in \mathbb{N}, \quad \exists K_j, \, r_j > 0, \quad \forall x \in \mathbb{R}, \qquad |f^{(j)}(x)| \le K_j (1 + |x|)^{r_j} e^{-|x|}.$$

The first system $(\Omega_{1,0})$ has been solved in [10, Proposition 2.3 and Lemma 3.1]. We recall

$$(\Omega_{1,0}) \begin{cases} \mathcal{L}A_{1,0} + a_{1,0}(3Q - 2Q^4) = 4Q^3 \\ (\mathcal{L}B_{1,0})' + a_{1,0}(3Q'') - 3A''_{1,0} - 4Q^3A_{1,0} = 4Q^3. \end{cases}$$

We found a solution $(a_{1,0}, A_{1,0}, B_{1,0})$ of $(\Omega_{1,0})$ such that $A_{1,0}$ is even, $B_{1,0}$ is odd, and

$$A_{1,0} = \frac{1}{3} (Q'(\int_0^x Q^2) - 2Q^3) + 2 \frac{\int Q}{\int Q^2} (-\frac{1}{3}Q - \frac{3}{2}xQ') \in \mathcal{Y},$$

$$B_{1,0} = \overline{B}_{1,0} + b_{1,0} (-\frac{Q'}{Q}), \quad \overline{B}_{1,0} \in \mathcal{Y},$$

$$a_{1,0} = -2 \frac{\int Q}{\int Q^2} < 0, \quad b_{1,0} = \lim_{x \to +\infty} B_{1,0} = -\frac{1}{2} \int Q^3 + \frac{1}{6} \frac{(\int Q)^2}{\int Q^2} < 0.$$
(5)

(Note that by explicit computations $\lim_{\pm\infty}(-\frac{Q'}{Q})=\pm 1$.) The interpretation of the resolution of the first system is the following: $a_{1,0}$ is related to the shift on the first soliton due to the collision and $b_{1,0}$ is related to the shift on the second soliton. In particular, in contrast with the cases of the nonlinearities u^2 and u^3 , the shifts on Q and Q_c are both negative, see [10, Remark 2]. Note that the explicit computation of $\overline{B}_{1,0}$ is possible but it is not required in [10].

For the system $(\Omega_{2,0})$, which writes as follows (see [10, Proposition 2.1]):

$$(\Omega_{2,0}) \begin{cases} (\mathcal{L}A_{2,0})' + a_{2,0}(3Q - 2Q^4)' = (6Q^2(1 + A_{1,0})^2)' \\ - a_{1,0}(4Q^3 + 3A_{1,0}'' + 4Q^3A_{1,0})' + 3a_{1,0}^2Q^{(3)}, \end{cases}$$

$$(\mathcal{L}B_{2,0})' + 3a_{2,0}Q'' - 3A_{2,0}'' - 4Q^3A_{2,0}$$

$$= 6Q^2(1 + A_{1,0})^2 + (6Q^2B_{1,0}(1 + A_{1,0}))' - \frac{1}{2}a_{1,0}(9A_{1,0}' + 3B_{1,0}'' + 4Q^3B_{1,0})' + \frac{3}{2}a_{1,0}^2Q'', \end{cases}$$

we claim the following, which is the main result of this note.

Theorem A. Let $(a_{1,0}, A_{1,0}, B_{1,0})$ satisfy $(\Omega_{1,0})$ and (5). Then, there exists a solution $(a_{2,0}, A_{2,0}, B_{2,0})$ of $(\Omega_{2,0})$ such that $A_{2,0}$ is even, $B_{2,0}$ is odd, and

$$A_{2,0} \in \mathcal{Y}, \quad B_{2,0} = \overline{B}_{2,0} + b_{2,0} \left(-\frac{Q'}{Q}\right), \quad \overline{B}_{2,0} \in \mathcal{Y},$$

$$b_{2,0} = \lim_{+\infty} B_{2,0} = -\frac{1}{18} \int Q^2 - \frac{3}{4} \frac{\left(\int Q\right)\left(\int Q^3\right)}{\int Q^2} - \frac{1}{18} \frac{\left(\int Q\right)^3}{\left(\int Q^2\right)^2} < 0.$$

The fact that $b_{2,0} \neq 0$ is fundamental in the proof of [10, Theorem 1]. Note that the explicit computation of $b_{2,0}$ in this note does not require the complete resolution of the system $(\Omega_{2,0})$. A main difficulty in solving systems $(\Omega_{k,\ell})$ is to find explicit antecedents by the operator \mathcal{L} . In the proof of Theorem A, we use the resolution of $(\Omega_{1,0})$, $b_{2,0} = \lim_{+\infty} B_{2,0} = \lim_{+\infty} \mathcal{L}B_{2,0}$ and the self-adjointness of \mathcal{L} to avoid computing $a_{2,0}$, $A_{2,0}$, and $B_{2,0}$. The structure of operator \mathcal{L} is described in Lemma 1.1 in section 1.

Formally, the fact that $b_{2,0} \neq 0$ means that for $t = T_c = c^{-\frac{1}{2} - \frac{1}{100}}$, (i.e., long after the collision time t = 0), we have

$$v(T_c, x) - Q(y) \sim Q_c(y_c) - b_{1,0}Q'_c(y_c) - b_{2,0}(Q_c^2)'(y_c)$$

 $\sim Q_c(y_c - b_{1,0}) - b_{2,0}(Q_c^2)'(y_c).$

Thus, $v(T_c)$ does not match a two soliton solution because of the nonzero term $b_{2,0}(Q_c^2)'(y_c)$ (for details, see proof of [10, Proposition 3.1]). Let u(t,x) be the asymptotic 2-soliton solution at $-\infty$, with speeds 1, c. Then, it follows from the analysis in [10] that the size of $(b_{2,0}Q_c^2)'$ measures the distance of u(t,x) from a pure two soliton solution after the collision (some analysis arguments are required in addition to

deal with the asymptotics in large time, see [9]). Recall that this is the first rigorous description of inelastic collision of two solitons.

The plan of this note is as follows. In section 1, we prove Theorem A. In section 2, we present some similar computations for the case u^2 in equation (1), used in [10]. The appendix contains the proof of some technical results.

1. Proof of Theorem A

We first recall the following elementary properties of the operator \mathcal{L} (see Lemma 2.2 in [10]).

Lemma 1.1 (Properties of \mathcal{L} , [10]). The operator \mathcal{L} defined in $L^2(\mathbb{R})$ by

$$\mathcal{L}f = -f'' + f - 4Q^3f$$

is self-adjoint and satisfies the following properties:

- (i) First eigenfunction: $\mathcal{L}Q^{\frac{5}{2}} = -\frac{21}{4}Q^{\frac{5}{2}}$.
- (ii) Second eigenfunction: $\mathcal{L}Q' = 0$; the kernel of \mathcal{L} is $\{\lambda Q', \lambda \in \mathbb{R}\}$.
- (iii) For any function $h \in L^2(\mathbb{R})$ orthogonal to Q' for the L^2 scalar product, there exists a unique function $f \in H^2(\mathbb{R})$ orthogonal to Q' such that $\mathcal{L}f = h$; moreover, if h is even (respectively, odd), then f is even (respectively, odd).
- (iv) Suppose that $f \in H^2(\mathbb{R})$ is such that $\mathcal{L}f \in \mathcal{Y}$, then $f \in \mathcal{Y}$.

Proof of Theorem A. The existence part, with the required structure on $A_{2,0}$, $B_{2,0}$, was proved in [10, Proposition 2.3]. In this proof, for the sake of simplicity, we denote

$$a_{1,0} = a_1$$
, $A_{1,0} = A_1$, $B_{1,0} = B_1$, $a_{2,0} = a_2$, $A_{2,0} = A_2$, $B_{2,0} = B_2$.

From [10, Proposition 2.1], system (Ω_2) writes

$$(\Omega_2) \begin{cases} (\mathcal{L}A_2)' + a_2(3Q - 2Q^4)' = (6Q^2(1 + A_1)^2)' \\ -a_1(4Q^3 + 3A_1'' + 4Q^3A_1)' + 3a_1^2Q^{(3)}, \\ (\mathcal{L}B_2)' + 3a_2Q'' - 3A_2'' - 4Q^3A_2 \\ = 6Q^2(1 + A_1)^2 + (6Q^2B_1(1 + A_1))' \\ -\frac{1}{2}a_1(9A_1' + 3B_1'' + 4Q^3B_1)' + \frac{3}{2}a_1^2Q'', \end{cases}$$

where (a_1, A_1, B_1) satisfy (Ω_1) and (5). Recall also the following functions used in the resolution of the system (Ω_1) in [10, Lemma 3.1]:

$$V_0 = -\frac{1}{3}Q - \frac{3}{2}xQ', \quad \mathcal{L}V_0 = 3Q - 2Q^4,$$

$$V_1 = \frac{1}{3}\left(Q'\left(\int_0^x Q^2\right) - 2Q^3\right), \quad \mathcal{L}V_1 = 4Q^3,$$

$$Z_0 = 3Q'' + 3V_0'' + 4Q^3V_0, \quad Z_1 = 3V_1'' + 4Q^3V_1 + 4Q^3.$$

We now claim

$$2b_2 = -a_1 \left(6 \int Q^2 (1 + A_1)^2 V_0 - a_1 \int (\mathcal{L}B_1)' V_0 - \frac{9}{2} a_1 \int A_1 Q'' + \frac{3}{2} a_1^2 \int Q'' Q + \frac{1}{2} \int A_1 (\mathcal{L}B_1)' \right) + 6 \int Q^2 (1 + A_1)^2 (1 + V_1) - a_1 \int (\mathcal{L}B_1)' V_1, \quad (6)$$

Proof of (6). Denote by V_2 the (even) solution of

$$\mathcal{L}V_2 = 6Q^2(1+A_1)^2 - a_1(4Q^3 + 3A_1'' + 4Q^3A_1) + 3a_1^2Q''.$$

Then, $A_2 = V_2 - a_2 V_0$ is solution of the first line of the system (Ω_2) . Moreover, since $(\mathcal{L}B_1)' = 4Q^3 - 3a_1Q'' + 3A_1'' + 4Q^3A_1$,

$$\mathcal{L}V_2 = 6Q^2(1+A_1)^2 - a_1(\mathcal{L}B_1)'. \tag{7}$$

Recall from [10, Proposition 2.3] that the functions A_2 , B_2 have the following structure: $A_2 \in \mathcal{Y}$ is even and $B_2 = \overline{B}_2 + b_2 \varphi$ is odd, where $\overline{B}_2 \in \mathcal{Y}$ and $\lim_{\pm \infty} \varphi = \pm 1$, $\varphi' \in \mathcal{Y}$ ($\varphi = -\frac{Q'}{Q}$, see [10, Claim 2.1]). Thus,

$$2b_2 = \lim_{+\infty} B_{2,0} - \lim_{-\infty} B_{2,0} = \lim_{+\infty} \mathcal{L}B_{2,0} - \lim_{-\infty} \mathcal{L}B_{2,0} = \int (\mathcal{L}B_2)'.$$

Therefore, from the second line of (Ω_2) , and then by (7), we obtain

$$2b_2 = 4 \int Q^3 A_2 + 6 \int Q^2 (1+A_1)^2 = 4 \int Q^3 V_2 - 4a_2 \int Q^3 V_0 + 6 \int Q^2 (1+A_1)^2$$
$$= \int (\mathcal{L}V_2)V_1 - 4a_2 \int Q^3 V_0 + 6 \int Q^2 (1+A_1)^2$$
$$= 6 \int Q^2 (1+A_1)^2 (1+V_1) - a_1 \int (\mathcal{L}B_1)' V_1 - 4a_2 \int Q^3 V_0.$$

Since $V_0 = -\frac{1}{3}Q - \frac{3}{2}xQ'$ and $\int Q^4 = \int Q$, we have

$$-4\int Q^3V_0 = \frac{4}{3}\int Q^4 + 6\int xQ'Q^3 = -\frac{1}{6}\int Q = \frac{1}{12}a_1\int Q^2.$$

Thus,

$$2b_2 = \left(\frac{1}{12} \int Q^2\right) a_1 a_2 + 6 \int Q^2 (1 + A_1)^2 (1 + V_1) - a_1 \int (\mathcal{L}B_1)' V_1. \tag{8}$$

Now, we determine a_2 from the equation of B_2 , where we have replaced $A_2 = V_2 - a_2 V_0$,

$$\begin{split} (\mathcal{L}B_2)' + a_2 Z_0 &= 3V_2'' + 4Q^3 V_2 + 6Q^2 (1 + A_1)^2 - \frac{9}{2} a_1 A_1'' + \frac{3}{2} a_1^2 Q'' \\ &- \frac{1}{2} a_1 (3B_1'' + 4Q^3 B_1)' + 6(Q^2 (1 + A_1) B_1)'. \end{split}$$

Multiplying by Q, integrating, and using $\mathcal{L}Q' = 0$ and

$$\int Z_0 Q = -\frac{1}{12} \int Q^2 \tag{9}$$

(see proof of (9) at the end of appendix A), we obtain

$$-\frac{1}{12}\left(\int Q^2\right)a_2$$

$$=\int (3V_2'' + 4Q^3V_2)Q + 6\int Q^3(1+A_1)^2 - \frac{9}{2}a_1\int A_1''Q + \frac{3}{2}a_1^2\int Q''Q$$

$$+\int \left[\frac{1}{2}a_1(3B_1'' + 4Q^3B_1) - 6Q^2B_1 - 6Q^2A_1B_1\right]Q'$$

$$=\int V_2(3Q'' + 4Q^4) + 6\int Q^3(1+A_1)^2 - \frac{9}{2}a_1\int A_1''Q + \frac{3}{2}a_1^2\int Q''Q$$

$$+\int \left[\frac{1}{2}a_1(3Q''' + 4Q^3Q') - 6Q^2Q' - 6Q^2A_1Q'\right]B_1.$$

On the one hand,

$$3Q'' + 4Q^4 = 3Q + Q^4 = \mathcal{L}(V_0 - Q)$$

since $\mathcal{L}Q = -3Q^4$. On the other hand, since $\mathcal{L}A_1 = 4Q^3 - a_1(3Q'' + Q^4)$,

$$\frac{1}{2}a_1(3Q''' + 4Q^3Q') - 6Q^2Q' - 6Q^2A_1Q' = -\frac{1}{2}(\mathcal{L}A_1)' - \frac{1}{2}A_1(4Q^3)'
= -\frac{1}{2}((\mathcal{L}A_1)' + 4(Q^3)'A_1) = -\frac{1}{2}\mathcal{L}(A_1').$$

Thus, by (7), since $\mathcal{L}Q' = 0$,

$$-\frac{1}{12} \left(\int Q^2 \right) a_2 = \int (\mathcal{L}V_2)(V_0 - Q) + 6 \int Q^3 (1 + A_1)^2$$

$$-\frac{9}{2} a_1 \int A_1'' Q + \frac{3}{2} a_1^2 \int Q'' Q + \frac{1}{2} \int A_1 (\mathcal{L}B_1)'$$

$$= 6 \int Q^2 (1 + A_1)^2 V_0 - a_1 \int (\mathcal{L}B_1)' (V_0 - Q)$$

$$-\frac{9}{2} a_1 \int A_1'' Q + \frac{3}{2} a_1^2 \int Q'' Q + \frac{1}{2} \int A_1 (\mathcal{L}B_1)'$$

$$= 6 \int Q^2 (1 + A_1)^2 V_0 - a_1 \int (\mathcal{L}B_1)' V_0$$

$$-\frac{9}{2} a_1 \int A_1'' Q + \frac{3}{2} a_1^2 \int Q'' Q + \frac{1}{2} \int A_1 (\mathcal{L}B_1)'.$$

By inserting this value of a_2 into (8), we obtain (6).

As $A_1 = V_1 - a_1 V_0$ and then $(\mathcal{L}B_1)' = 4Q^3 - 3a_1 Q'' + 3A_1'' + 4Q^3 A_1$, we obtain

$$2b_2 = 6 \int Q^2 (1 + A_1)^3 - \frac{3}{2} a_1 \int A_1 (\mathcal{L}B_1)' + \frac{9}{2} a_1^2 \int A_1 Q'' - \frac{3}{2} a_1^3 \int Q'' Q$$

$$= 6 \int Q^2 (1 + A_1)^3 - \frac{3}{2} a_1 \int A_1 (4Q^3 - 3a_1 Q'' + 3A_1'' + 4Q^3 A_1)$$

$$+ \frac{9}{2} a_1^2 \int A_1 Q'' - \frac{3}{2} a_1^3 \int Q'' Q$$

$$= 6 \int Q^2 (1 + A_1)^3 - \frac{3}{2} a_1 \int A_1 (4Q^3 - 6a_1 Q'' + 3A_1'' + 4Q^3 A_1) - \frac{3}{2} a_1^3 \int Q'' Q.$$

Since

$$\mathcal{L}A_1 + a_1(3Q - 2Q^4) = -A_1'' + A_1 - 4Q^3A_1 + 3a_1Q - 2a_1Q^4 = 4Q^3$$

we deduce

$$4Q^3 - 6a_1Q'' + 3A_1'' + 4Q^3A_1 = -8Q^3(1+A_1) + 3(A_1 + a_1Q),$$

and so

$$2b_2 = 6 \int Q^2 (1+A_1)^3 + \frac{3}{2}a_1 \int A_1 (8Q^3 (1+A_1) - 3(A_1 + a_1 Q)) - \frac{3}{2}a_1^3 \int Q'' Q.$$

We use again $A_1 = V_1 - a_1 V_0$ and we sort terms by increasing powers of a_1 . We get

$$2b_2 = 6 \int Q^2 (1 + V_1 - a_1 V_0)^3 + 12a_1 \int (V_1 - a_1 V_0) Q^3 (1 + V_1 - a_1 V_0)$$
$$- \frac{9}{2} a_1 \int (V_1 - a_1 V_0) (V_1 + a_1 (-V_0 + Q)) - \frac{3}{2} a_1^3 \int Q'' Q$$
$$= \mu_0 + a_1 \mu_1 + a_1^2 \mu_2 + a_1^3 \mu_3,$$

where

$$\begin{split} &\mu_0 = 6 \int Q^2 (1+V_1)^3, \\ &\mu_1 = -18 \int Q^2 (1+V_1)^2 V_0 + 12 \int V_1 (1+V_1) Q^3 - \frac{9}{2} \int V_1^2, \\ &\mu_2 = 18 \int Q^2 (1+V_1) V_0^2 - 12 \int V_0 (1+2V_1) Q^3 + \frac{9}{2} \int V_1 (2V_0 - Q), \\ &\mu_3 = -6 \int Q^2 V_0^3 + 12 \int Q^3 V_0^2 - \frac{9}{2} \int V_0^2 + \frac{9}{2} \int V_0 Q - \frac{3}{2} \int Q'' Q. \end{split}$$

We claim the following values for μ_0 , μ_1 , μ_2 , and μ_3 .

Lemma 1.2.

$$\mu_0 = -\frac{1}{9} \int Q^2, \quad \mu_1 = \frac{3}{4} \int Q^3, \quad \mu_2 = \frac{1}{24} \int Q, \quad \mu_3 = \frac{5}{144} \int Q^2.$$

Assuming Lemma 1.2, and using $a_1 = -2 \int Q/\int Q^2$, we obtain

$$2b_2 = -\frac{1}{9} \int Q^2 - \frac{3}{2} \frac{(\int Q)(\int Q^3)}{\int Q^2} - \frac{1}{9} \frac{(\int Q)^3}{(\int Q^2)^2}.$$

This completes the proof of Theorem A provided we prove Lemma 1.2.

Proof of Lemma 1.2. The main ingredient is Lemma A.1 (see appendix).

• Computation of μ_0 .

$$\mu_0 = 6 \int Q^2 + 18 \int Q^2 V_1 + 18 \int Q^2 V_1^2 + 6 \int Q^2 V_1^3.$$

By Lemma A.1,

$$18 \int Q^{2}V_{1} = 6 \int Q^{2}(Q'(\int_{0}^{x} Q^{2}) - 2Q^{3})$$

$$= -2 \int Q^{5} - 12 \int Q^{5} = -14 \int Q^{5} = -20 \int Q^{2},$$

$$18 \int Q^{2}V_{1}^{2} = 2 \int Q^{2}(Q'(\int_{0}^{x} Q^{2}) - 2Q^{3})^{2}$$

$$= 2 \int Q^{2}(Q')^{2}(\int_{0}^{x} Q^{2})^{2} - 8 \int Q^{5}Q'(\int_{0}^{x} Q^{2}) + 8 \int Q^{8}$$

and

$$\begin{split} 6 \int Q^2 V_1^3 &= \frac{2}{9} \int Q^2 \big(Q' \big(\int_0^x Q^2 \big) - 2 Q^3 \big)^3 \\ &= \frac{2}{9} \int Q^2 (Q')^3 \big(\int_0^x Q^2 \big)^3 - \frac{4}{3} \int Q^5 (Q')^2 \big(\int_0^x Q^2 \big)^2 \\ &\quad + \frac{8}{3} \int Q^8 Q' \big(\int_0^x Q^2 \big) - \frac{16}{9} \int Q^{11} \\ &= \frac{2}{9} \int Q^4 Q' \big(\int_0^x Q^2 \big)^3 - \frac{4}{45} \int Q^7 Q' \big(\int_0^x Q^2 \big)^3 - \frac{4}{3} \int Q^7 \big(\int_0^x Q^2 \big)^2 \\ &\quad + \frac{8}{15} \int Q^{10} \big(\int_0^x Q^2 \big)^2 - \frac{56}{27} \int Q^{11} \\ &= -\frac{22}{15} \int Q^7 \big(\int_0^x Q^2 \big)^2 + \frac{17}{30} \int Q^{10} \big(\int_0^x Q^2 \big)^2 - \frac{56}{11} \int Q^{11}. \end{split}$$

 $=2\int Q^4 \left(\int_0^x Q^2\right)^2 -\frac{4}{5}\int Q^7 \left(\int_0^x Q^2\right)^2 +\frac{28}{3}\int Q^8,$

By Lemma A.1,
$$\int Q^{10} \left(\int_0^x Q^2 \right)^2 = \frac{35}{17} \int Q^7 \left(\int_0^x Q^2 \right)^2 - \frac{10}{9 \cdot 17} \int Q^{11}$$
. Thus,
$$6 \int Q^2 V_1^3 = -\frac{3}{10} \int Q^7 \left(\int_0^x Q^2 \right)^2 - \frac{19}{9} \int Q^{11}$$
$$= -\frac{3}{10} \int Q^7 \left(\int_0^x Q^2 \right)^2 - \frac{40}{9} \int Q^8.$$

Therefore,

$$18 \int Q^2 V_1^2 + 6 \int Q^2 V_1^3 = 2 \int Q^4 \left(\int_0^x Q^2 \right)^2 - \frac{11}{10} \int Q^7 \left(\int_0^x Q^2 \right)^2 + \frac{44}{9} \int Q^8.$$

By Lemma A.1, $\int Q^7 \left(\int_0^x Q^2 \right)^2 = \frac{20}{11} \int Q^4 \left(\int_0^x Q^2 \right)^2 - \frac{5}{33} \int Q^8$, and so

$$18 \int Q^2 V_1^2 + 6 \int Q^2 V_1^3 = \frac{91}{18} \int Q^8.$$

Finally,

$$\mu_0 = -14 \int Q^2 + \frac{91}{18} \int Q^8 = -\frac{1}{9} \int Q^2,$$

since, by Lemma A.1, $\int Q^8 = \frac{25}{13} \int Q^5 = \frac{250}{7 \cdot 13} \int Q^2$.

• Computation of μ_1 .

$$\mu_1 = -18 \int Q^2 (1+V_1)^2 V_0 + 12 \int Q^3 V_1 (1+V_1) - \frac{9}{2} \int V_1^2.$$

First.

$$-18 \int Q^2 (1+V_1)^2 V_0 = -18Q^2 V_0 - 36 \int Q^2 V_1 V_0 - 18 \int Q^2 V_1^2 V_0.$$

We use that $V_0 = -\frac{1}{3}Q - \frac{3}{2}xQ'$:

$$-18Q^{2}V_{0} = 6 \int Q^{3} + 27 \int xQ'Q^{2} = -3 \int Q^{3},$$

$$-36 \int Q^{2}V_{1}V_{0} = 4 \int Q^{2}(Q'(\int_{0}^{x}Q^{2}) - 2Q^{3})(Q + \frac{9}{2}xQ')$$

$$= 4 \int Q^{3}Q'(\int_{0}^{x}Q^{2}) + 18 \int Q^{2}(Q')^{2}x(\int_{0}^{x}Q^{2})$$

$$-8 \int Q^{6} - 36 \int xQ'Q^{5}$$

$$= -3 \int Q^{6} + 18 \int xQ^{4}(\int_{0}^{x}Q^{2}) - \frac{36}{5} \int xQ^{7}(\int_{0}^{x}Q^{2}),$$

and

$$-18 \int Q^{2}V_{1}^{2}V_{0} = \frac{2}{3} \int Q^{2}(Q'(\int_{0}^{x}Q^{2}) - 2Q^{3})^{2}(Q + \frac{9}{2}xQ')$$

$$= \frac{2}{3} \int (Q^{3} + \frac{9}{2}xQ'Q^{2})$$

$$\times (Q^{2}(\int_{0}^{x}Q^{2})^{2} - \frac{2}{5}(\int_{0}^{x}Q^{2})^{2} - 4Q^{3}Q'(\int_{0}^{x}Q^{2}) + 4Q^{6})$$

$$= \frac{2}{3} \int Q^{5}(\int_{0}^{x}Q^{2})^{2} - \frac{4}{15} \int Q^{8}(\int_{0}^{x}Q^{2})^{2} - \frac{8}{3} \int Q^{6}Q'(\int_{0}^{x}Q^{2})$$

$$+ \frac{8}{3} \int Q^{9} + 3 \int xQ'Q^{4}(\int_{0}^{x}Q^{2})^{2} - \frac{6}{5} \int xQ'Q^{7}(\int_{0}^{x}Q^{2})^{2}$$

$$- 12 \int xQ^{7}(\int_{0}^{x}Q^{2}) + \frac{24}{5} \int xQ^{10}(\int_{0}^{x}Q^{2}) + 12 \int xQ'Q^{8}$$

$$= \frac{1}{15} \int Q^{5}(\int_{0}^{x}Q^{2})^{2} - \frac{7}{60} \int Q^{8}(\int_{0}^{x}Q^{2})^{2} + \frac{12}{7} \int Q^{9}$$

$$- \frac{66}{5} \int xQ^{7}(\int_{0}^{x}Q^{2}) + \frac{51}{10} \int xQ^{10}(\int_{0}^{x}Q^{2}).$$

Summing up

$$-18 \int Q^{2} (1+V_{1})^{2} V_{0} = -3 \int Q^{3} - 3 \int Q^{6} + \frac{12}{7} \int Q^{9} + 18 \int x Q^{4} (\int_{0}^{x} Q^{2})$$
$$- \frac{102}{5} \int x Q^{7} (\int_{0}^{x} Q^{2}) + \frac{51}{10} \int x Q^{10} (\int_{0}^{x} Q^{2})$$
$$+ \frac{1}{15} \int Q^{5} (\int_{0}^{x} Q^{2})^{2} - \frac{7}{60} \int Q^{8} (\int_{0}^{x} Q^{2})^{2}.$$

By Lemma A.1,

$$\int xQ^{10} \left(\int_0^x Q^2 \right) = \frac{35}{17} \int xQ^7 \left(\int_0^x Q^2 \right) - \frac{80}{7 \cdot 9 \cdot 17} \int Q^9,$$

thus

$$-18 \int Q^{2} (1+V_{1})^{2} V_{0} = -3 \int Q^{3} - 3 \int Q^{6} + \frac{4}{3} \int Q^{9} + \frac{1}{15} \int Q^{5} (\int_{0}^{x} Q^{2})^{2}$$
$$-\frac{7}{60} \int Q^{8} (\int_{0}^{x} Q^{2})^{2} + 18 \int x Q^{4} (\int_{0}^{x} Q^{2})$$
$$-\frac{99}{10} \int x Q^{7} (\int_{0}^{x} Q^{2}).$$

From Lemma A.1,

so that

$$-18 \int Q^{2} (1+V_{1})^{2} V_{0} = -3 \int Q^{3} - \frac{9}{8} \int Q^{6} + \frac{4}{3} \int Q^{9} + \frac{1}{15} \int Q^{5} (\int_{0}^{x} Q^{2})^{2} - \frac{7}{60} \int Q^{8} (\int_{0}^{x} Q^{2})^{2},$$

and

$$\begin{aligned} 12 \int V_1(1+V_1)Q^3 &= 12 \int V_1Q^3 + 12 \int V_1^2Q^3 \\ &= 4 \int (Q'(\int_0^x Q^2) - 2Q^3)Q^3 \\ &\quad + \frac{4}{3} \int (Q^2(\int_0^x Q^2)^2 - \frac{2}{5}(\int_0^x Q^2)^2 - 4Q'Q^3(\int_0^x Q^2) + 4Q^6)Q^3 \\ &= -9 \int Q^6 + \frac{128}{21} \int Q^9 + \frac{4}{3} \int Q^5(\int_0^x Q^2)^2 \\ &\quad - \frac{8}{15} \int Q^8(\int_0^x Q^2)^2. \end{aligned}$$

Therefore,

$$-18 \int Q^{2} (1+V_{1})^{2} V_{0} + 12 \int V_{1} (1+V_{1}) Q^{3} = -3 \int Q^{3} - \frac{81}{8} \int Q^{6}$$

$$+ \frac{156}{21} \int Q^{9} + \frac{7}{5} \int Q^{5} (\int_{0}^{x} Q^{2})^{2}$$

$$- \frac{39}{60} \int Q^{8} (\int_{0}^{x} Q^{2})^{2}.$$

We use that

$$\int Q^8 \left(\int_0^x Q^2 \right)^2 = \frac{25}{13} \int Q^5 \left(\int_0^x Q^2 \right)^2 - \frac{10}{7 \cdot 13} \int Q^9.$$

Thus,

$$-18 \int Q^{2} (1+V_{1})^{2} V_{0} + 12 \int V_{1} (1+V_{1}) Q^{3} = -3 \int Q^{3} - \frac{81}{8} \int Q^{6} + \frac{15}{2} \int Q^{9} + \frac{3}{20} \int Q^{5} (\int_{0}^{x} Q^{2})^{2}$$
$$= -3 \int Q^{3} + \frac{39}{8} \int Q^{6} + \frac{3}{20} \int Q^{5} (\int_{0}^{x} Q^{2})^{2},$$

since $\int Q^9 = 2 \int Q^6$ (Lemma A.1).

$$\begin{split} -\frac{9}{2} \int V_1^2 &= -\frac{1}{2} \int \left(Q' \left(\int_0^x Q^2 \right)^2 - 2Q^3 \right)^2 \\ &= -\frac{1}{2} \int Q^2 \left(\int_0^x Q^2 \right)^2 + \frac{1}{5} \int Q^5 \left(\int_0^x Q^2 \right)^2 + 2 \int Q' Q^3 \left(\int_0^x Q^2 \right) - 2 \int Q^6 \\ &= -\frac{1}{2} \int Q^2 \left(\int_0^x Q^2 \right)^2 + \frac{1}{5} \int Q^5 \left(\int_0^x Q^2 \right)^2 - \frac{5}{2} \int Q^6. \end{split}$$

Summing up, we get

$$\mu_1 = -3 \int Q^3 + \frac{19}{6} \int Q^6 - \frac{1}{2} \int Q^2 \left(\int_0^x Q^2 \right)^2 + \frac{7}{20} \int Q^5 \left(\int_0^x Q^2 \right)^2,$$

but

$$\int Q^5 \left(\int_0^x Q^2 \right)^2 = \frac{10}{7} \int Q^2 \left(\int_0^x Q^2 \right)^2 - \frac{5}{14} \int Q^6,$$

and $\int Q^6 = \frac{5}{3} \int Q^3$, by Lemma A.1, so that

$$\mu_1 = \frac{3}{4} \int Q^3.$$

• Computation of μ_2 .

$$\mu_2 = 18 \int Q^2 (1 + V_1) V_0^2 - 12 \int Q^3 V_0 (1 + 2V_1) + \frac{9}{2} \int V_1 (2V_0 - Q).$$

First

$$18 \int Q^{2}(1+V_{1})V_{0}^{2} = \frac{2}{3} \int Q^{2}(3+Q'(\int_{0}^{x}Q^{2})-2Q^{3})(Q+\frac{9}{2}xQ')^{2}$$

$$= \frac{2}{3} \int (Q^{4}+9xQ^{3}Q'+\frac{81}{4}x^{2}Q^{4}-\frac{81}{10}x^{2}Q^{7})$$

$$\times (3+Q'(\int_{0}^{x}Q^{2})-2Q^{3})$$

$$= 2 \int Q^{4}+18 \int xQ^{3}Q'+\frac{81}{2} \int x^{2}Q^{4}$$

$$-\frac{81}{4} \int x^{2}Q^{7}+\frac{2}{3} \int Q^{4}Q'(\int_{0}^{x}Q^{2})$$

$$+6 \int xQ^{5}(\int_{0}^{x}Q^{2})-\frac{12}{5} \int xQ^{8}(\int_{0}^{x}Q^{2})$$

$$+\frac{27}{2} \int x^{2}Q^{4}Q'(\int_{0}^{x}Q^{2})-\frac{27}{5} \int x^{2}Q^{7}Q'(\int_{0}^{x}Q^{2})-\frac{4}{3} \int Q^{7}$$

$$-12 \int xQ^{6}Q'-27 \int x^{2}Q^{7}+\frac{54}{5} \int x^{2}Q^{10}.$$

Thus,

$$18 \int Q^{2}(1+V_{1})V_{0}^{2} = -\frac{5}{2} \int Q^{4} + \left(-\frac{22}{15} + \frac{12}{7}\right) \int Q^{7} + \frac{81}{2} \int x^{2}Q^{4} - \frac{459}{10} \int x^{2}Q^{7} + \frac{17 \cdot 27}{40} \int x^{2}Q^{10} + \frac{3}{5} \int xQ^{5}(\int_{0}^{x}Q^{2}) - \frac{21}{20} \int xQ^{8}(\int_{0}^{x}Q^{2}).$$

But (Lemma A.1)

$$\int x^2 Q^{10} = \frac{35}{17} \int x^2 Q^7 - \frac{10}{7 \cdot 17} \int Q^7,$$

and

$$\int x^2 Q^7 = \frac{20}{11} \int x^2 Q^4 - \frac{5}{22} \int Q^4,$$

so that

$$-\frac{459}{10} \int x^2 Q^7 + \frac{17 \cdot 27}{40} \int x^2 Q^{10} = -\frac{81}{2} \int x^2 Q^4 + \frac{81}{16} \int Q^4 - \frac{27}{28} \int Q^7.$$

Therefore.

$$18 \int Q^{2}(1+V_{1})V_{0}^{2} = \frac{41}{16} \int Q^{4} - \frac{43}{60} \int Q^{7}$$

$$+ \frac{3}{5} \int xQ^{5} \left(\int_{0}^{x} Q^{2}\right) - \frac{21}{20} \int xQ^{8} \left(\int_{0}^{x} Q^{2}\right),$$

$$-12 \int V_{0}(1+2V_{1})Q^{3} = 4 \int \left(Q + \frac{9}{2}xQ'\right)Q^{3} \left(1 + \frac{2}{3}Q'\left(\int_{0}^{x} Q^{2}\right) - \frac{4}{3}Q^{3}\right)$$

$$= 4 \int \left(Q^{4} + \frac{9}{2}xQ'Q^{3}\right) + \frac{8}{3} \int Q^{4}Q'\left(\int_{0}^{x} Q^{2}\right)$$

$$+ 12 \int xQ^{5} \left(\int_{0}^{x} Q^{2}\right) - \frac{24}{5} \int xQ^{8} \left(\int_{0}^{x} Q^{2}\right) - \frac{16}{3} \int Q^{7}$$

$$- 24 \int xQ^{6}Q'$$

$$= -\frac{1}{2} \int Q^{4} - \frac{8}{15} \int Q^{7} + 12 \int xQ^{5} \left(\int_{0}^{x} Q^{2}\right)$$

$$- \frac{24}{5} \int xQ^{8} \left(\int_{0}^{x} Q^{2}\right) - \frac{40}{21} \int Q^{7}.$$

Thus.

$$18 \int Q^{2}(1+V_{1})V_{0}^{2} - 12 \int V_{0}(1+2V_{1})Q^{3}$$

$$= \frac{33}{16} \int Q^{4} - \frac{265}{84} \int Q^{7} + \frac{63}{5} \int xQ^{5}(\int_{0}^{x} Q^{2}) - \frac{117}{20} \int xQ^{8}(\int_{0}^{x} Q^{2}).$$

But, by Lemma A.1,

$$\int xQ^8 \left(\int_0^x Q^2 \right) = \frac{25}{13} \int xQ^5 \left(\int_0^x Q^2 \right) - \frac{12}{91} \int Q^7,$$

and so

$$18 \int Q^{2}(1+V_{1})V_{0}^{2} - 12 \int V_{0}(1+2V_{1})Q^{3}$$

$$= \frac{33}{16} \int Q^{4} - \frac{143}{60} \int Q^{7} + \frac{27}{20} \int xQ^{5}(\int_{0}^{x} Q^{2}).$$

As $\int Q^7 = \frac{20}{11} \int Q^4$, we get

$$18 \int Q^2 (1+V_1)V_0^2 - 12 \int V_0 (1+2V_1)Q^3 = -\frac{109}{48} \int Q^4 + \frac{27}{20} \int xQ^5 \left(\int_0^x Q^2\right)$$

and

$$\frac{9}{2} \int V_1(2V_0 - Q) = -\frac{1}{2} \int (Q'(\int_0^x Q^2) - 2Q^3)(5Q + 9xQ')$$

$$= -\frac{5}{2} \int QQ'(\int_0^x Q^2) - \frac{9}{2} \int xQ^2(\int_0^x Q^2) + \frac{9}{5} \int xQ^5(\int_0^x Q^2)$$

$$+ 5 \int Q^4 + 9 \int xQ^3Q'$$

$$= 4 \int Q^4 - \frac{9}{2} \int xQ^2(\int_0^x Q^2) + \frac{9}{5} \int xQ^5(\int_0^x Q^2).$$

Finally,

$$\mu_2 = \frac{83}{48} \int Q^4 + \frac{63}{20} \int xQ^5 \left(\int_0^x Q^2 \right) - \frac{9}{2} \int xQ^2 \left(\int_0^x Q^2 \right).$$

But, by Lemma A.1,

$$\frac{63}{20} \int x Q^5 \left(\int_0^x Q^2 \right) = \frac{9}{2} \int x Q^2 \left(\int_0^x Q^2 \right) - \frac{27}{16} \int Q^4.$$

Thus

$$\mu_2 = \frac{1}{24} \int Q^4 = \frac{1}{24} \int Q.$$

• Computation of μ_3 .

$$\mu_3 = -6 \int Q^2 V_0^3 + 12 \int Q^3 V_0^2 - \frac{9}{2} \int V_0^2 + \frac{9}{2} \int V_0 Q - \frac{3}{2} \int Q'' Q.$$

First.

$$\begin{split} -6\int Q^2V_0^3 + 12\int Q^3V_0^2 &= 6\int V_0^2(-V_0 + 2Q) \\ &= \frac{2}{9}\int (Q + \frac{9}{2}xQ')^2Q^2(7Q + \frac{9}{2}xQ') \\ &= \frac{2}{9}\int (Q^2 + 9xQ'Q + \frac{81}{4}x^2Q^2 - \frac{81}{10}x^2Q^5) \\ &\quad \times (7Q^3 + \frac{9}{2}xQ^2Q') \\ &= \frac{14}{9}\int Q^5 + \int xQ^4Q' + 14\int xQ^4Q' + 9\int x^2Q^5 \\ &\quad - \frac{18}{5}\int x^2Q^8 + \frac{63}{2}\int x^2Q^5 + \frac{81}{4}\int x^3Q^4Q' \\ &\quad - \frac{63}{5}\int x^2Q^8 - \frac{81}{10}\int x^3Q^7Q' \\ &= -\frac{13}{9}\int Q^5 + \frac{567}{20}\int x^2Q^5 - \frac{1053}{80}\int x^2Q^8. \end{split}$$

But, by Lemma A.1,

$$\int x^2 Q^8 = \frac{25}{13} \int x^2 Q^5 - \frac{2}{13} \int Q^5,$$

and so

$$-6\int Q^2V_0^3 + 12\int Q^3V_0^2 = \frac{209}{360}\int Q^5 + \frac{243}{80}\int x^2Q^5.$$

Second,

$$\frac{9}{2} \int V_0(-V_0 + Q) = -\frac{1}{2} \int (Q + \frac{9}{2}xQ')(4Q + \frac{9}{2}xQ')$$

$$= -\frac{1}{2} \int (4Q^2 + \frac{9}{2}xQQ' + 18xQ'Q + \frac{81}{4}x^2Q^2 - \frac{81}{10}x^2Q^5)$$

$$= \frac{29}{8} \int Q^2 - \frac{81}{8} \int x^2Q^2 + \frac{81}{20} \int x^2Q^5.$$

Thus

$$\mu_3 = \frac{29}{8} \int Q^2 + \frac{209}{360} \int Q^5 - \frac{81}{8} \int x^2 Q^2 + \frac{567}{80} \int x^2 Q^5.$$

But

$$\int x^2 Q^5 = \frac{10}{7} \int x^2 Q^2 - \frac{5}{7} \int Q^2 \quad \text{and} \quad \int Q^5 = \frac{10}{7} \int Q^2,$$

so that

$$\mu_3 = \frac{5}{144} \int Q^2.$$

2. Resolution of systems $(\Omega_{2,0})$ and $(\Omega_{1,1})$ for p=2

In this section, we consider the equation

$$\partial_t u + \partial_x (\partial_x^2 u + u^2) = 0, \quad x, t \in \mathbb{R}.$$

The problem is similar, except that there exist explicit multi-solitons which describe the collision of several solitons. Our objective in [10] was to point out that at the main orders our decomposition corresponds to the explicit 2-soliton solution. In doing so, we needed to compute the solutions of $(\Omega_{1,0})$, $(\Omega_{1,1})$ and $(\Omega_{2,0})$ (for the nonlinearity u^2 , the system $(\Omega_{1,1})$ appears at the same order of size as $(\Omega_{2,0})$). Note that the same could be done for u^3 .

In the case u^2 .

$$Q(x) = \frac{3}{2}\cosh^{-1}(\frac{x}{2}) \quad \text{solves} \quad Q'' + Q = Q^2.$$

We recall from [10, Lemma 3.1] that the resolution of $(\Omega_{1,0})$ gives

$$a_{1,0} = \frac{2}{3}, \quad A_{1,0} = -\frac{4}{3}Q, \quad B_{1,0} = 2\frac{Q'}{Q}.$$

We now solve $(\Omega_{2,0})$ and $(\Omega_{1,1})$.

Proposition 2.1 (Resolution of $(\Omega_{2,0})$ and $(\Omega_{1,1})$ for p=2).

$$a_{2,0}=-\frac{4}{9}, \ a_{1,1}=\frac{2}{3}, \ A_{2,0}=-2+\frac{4}{3}Q, \ A_{1,1}=2-\frac{2}{3}Q-\frac{1}{3}xQ', \ b_{2,0}=\lim_{+\infty}B_{2,0}=\frac{4}{3}.$$

for solutions to systems $(\Omega_{2,0})$ and $(\Omega_{1,1})$.

From [10, Corollary 3.1], there are several solutions. Indeed, in general, if $(a_{k,\ell}, A_{k,\ell}, B_{k,\ell})$ is solution of $(\Omega_{k,\ell})$, then for any $(\gamma_{k,\ell}, \delta_{k,\ell}) \in \mathbb{R}^2$,

$$(a_{k,\ell} + \gamma_{k,\ell} a_{1,0}, A_{k,\ell} + \gamma_{k,\ell} (1 + A_{1,0}), B_{k,\ell} + \gamma_{k,\ell} B_{1,0} + \delta_{k,\ell} Q')$$
(10)

is also solution, which gives two degrees of freedom at each step. The choice of the solution above is related to the exact 2-soliton solutions.

Proof of Proposition 2.1.

• System $(\Omega_{2,0})$.

From [10, Proposition 2.1],

$$\begin{cases}
(\mathcal{L}A_{2,0})' + a_{2,0}(3Q - 2Q^2)' = (-A_{1,0} + A_{1,0}^2)' - (3B_{1,0}'' + 2QB_{1,0}) \\
- a_{1,0}(Q + 3A_{1,0}'' + 2QA_{1,0})' + 3a_{1,0}^2Q^{(3)}, \\
(\mathcal{L}B_{2,0})' + 3a_{2,0}Q'' - 3A_{2,0}'' - 2QA_{2,0} \\
= A_{1,0} + A_{1,0}^2 + (-2B_{1,0} + A_{1,0}B_{1,0})' \\
- \frac{1}{2}a_{1,0}(9A_{1,0}' + 3B_{1,0}'' + 2QB_{1,0})' + \frac{3}{2}a_{1,0}^2Q''.
\end{cases}$$

Replacing $A_{1,0}$ and $B_{1,0}$, the first line of the system becomes

$$(\mathcal{L}A_{2,0})' = (\frac{4}{3}Q + \frac{16}{9}Q^2)' - (-2Q' + 4Q'QQ') - \frac{2}{3}(Q - 4Q'' - \frac{8}{3}Q^2)' + \frac{4}{3}Q^{(3)} - a_{2,0}(3Q - 2Q^2)' = (\frac{8}{3}Q - \frac{4}{9}Q^2 - a_{2,0}(3Q - 2Q^2))'.$$

We choose $\gamma_{2,0} = -2$ in (10), so that we consider $A_{2,0}$ solution of

$$\mathcal{L}A_{2,0} = \frac{8}{3}Q - \frac{4}{9}Q^2 - a_{2,0}(3Q - 2Q^2) - 2.$$

Consider the function V_2 defined by

$$V_2 = -2 + \frac{16}{9}Q + \frac{2}{3}xQ', \quad \mathcal{L}V_2 = \frac{8}{3}Q - \frac{4}{9}Q^2 - 2.$$

Then, $A_{2,0} = V_2 - a_{2,0}V_0$, where in this case $V_0 = -Q - \frac{3}{2}xQ'$ solves $\mathcal{L}V_0 = 3Q - 2Q^2$. Inserting the expressions of $a_{1,0}$, $A_{1,0}$, $B_{1,0}$, and $A_{2,0}$ into the second line of $(\Omega_{2,0})$ and after some calculations, we obtain the following equation for $B_{2,0}$:

$$(\mathcal{L}B_{2,0})' + a_{2,0}Z_0 = 3V_2'' + 2QV_2 + \frac{4}{3}Q + \frac{4}{9}Q^2,$$

where $Z_0 = 3Q'' + 3V_0'' + 2QV_0$. The value of $a_{2,0}$ is thus given by

$$a_{2,0} \int Z_0 Q = \int Z_2 Q,$$

where $Z_2 = 3V_2'' + 2QV_2 + \frac{4}{3}Q + \frac{4}{9}Q^2$. Recall that $\int Z_0Q = -\frac{3}{4}\int Q^2$ [10, Claim 2.3]. We claim

$$\int Z_2 Q = \frac{1}{3} \int Q. \tag{11}$$

Proof of (11).

$$\int Z_2 Q = \int V_2 (3Q'' + 2Q^2) + \frac{4}{3} \int Q^2 + \frac{4}{9} \int Q^3$$

$$= \int \mathcal{L}V_2 (-2Q - \frac{3}{2}xQ') + \frac{4}{3} \int Q^2 + \frac{4}{9} \int Q^3$$

$$= -\int (\frac{8}{3}Q - \frac{4}{9}Q^2 - 2)(2Q + \frac{3}{2}xQ') + \frac{4}{3} \int Q^2 + \frac{4}{9} \int Q^3$$

$$= \int Q - 2 \int Q^2 + \frac{10}{9} \int Q^3 = \frac{1}{3} \int Q.$$

Thus,

$$a_{2,0} = -\frac{4}{9}$$
 and $A_{2,0} = -2 + \frac{4}{3}Q$.

Finally, we compute $b_{2,0}$. As in section 1, we have $2b_{2,0} = \int (\mathcal{L}B_{2,0})'$. Thus, since $\int Q = 6$, $\int QV_0 = -\frac{1}{4} \int Q$, and $\int V_2Q = -\frac{5}{9} \int Q$, we obtain

$$2b_{2,0} = -2a_{2,0} \int V_0 Q + 2 \int V_2 Q + \frac{4}{3} \int Q + \frac{4}{9} \int Q^2 = \frac{4}{9} \int Q = \frac{8}{3}.$$

• System $(\Omega_{1,1})$.

The system $(\Omega_{1,1})$ writes

$$\begin{cases} (\mathcal{L}A_{1,1})' + a_{1,1}(3Q - 2Q^2)' = 3A'_{1,0} + 3B''_{1,0} + 2QB_{1,0} \\ (\mathcal{L}B_{1,1})' + 3a_{1,1}Q'' - 3A''_{1,1} - 2QA_{1,1} = 3B'_{1,0}. \end{cases}$$

Replacing $A_{1,0}$ and $B_{1,0}$ by their respective expressions, the first line of the system becomes

$$(\mathcal{L}A_{1,1})' = -2Q' - a_{1,1}(3Q - 2Q^2)'.$$

We choose $\gamma_{1,1} = 2$, so that

$$\mathcal{L}A_{1,1} = 2 - 2Q - a_{1,1}(3Q - 2Q^2)$$
 and $A_{1,1} = 2 - 2Q - xQ' - a_{1,1}V_0$.

Thus, the second line of the system becomes

$$(\mathcal{L}B_{1\,1})' + a_{1\,1}Z_0 = -2Q + 3(-2Q - xQ')'' + 2Q(2 - 2Q - xQ'),$$

and so

$$a_{1,1} \int Z_0 Q = -\frac{5}{2} \int Q^2 + \frac{5}{3} \int Q^3 = -\frac{1}{2} \int Q^2 = -\frac{1}{2} \int Q.$$

Since $\int Z_0 Q = -\frac{3}{4} \int Q$, we deduce $a_{1,1} = \frac{2}{3}$ and so $A_{1,1} = 2 - \frac{2}{3}Q - \frac{1}{3}xQ'$.

Appendix: Identities for solutions of (2)

Lemma A.1. Let Q be defined in (2). For any $r \geq 1$,

$$\int Q^{r+3} = \frac{5r}{2r+3} \int Q^r,\tag{12}$$

$$\int Q^{r+3} \left(\int_0^x Q^2 \right)^2 = \frac{5r}{2r+3} \int Q^r \left(\int_0^x Q^2 \right)^2 - \frac{10}{(2r+3)(r+2)} \int Q^{r+4}, \quad (13)$$

$$\int xQ^{r+3} \int_0^x Q^2 = \frac{5r}{2r+3} \int xQ^r \int_0^x Q^2 - \frac{10(r+1)}{(2r+3)r(r+2)} \int Q^{r+2}, \qquad (14)$$

$$\int x^2 Q^{r+3} = \frac{5r}{2r+3} \int x^2 Q^r - \frac{10}{r(2r+3)} \int Q^r.$$
 (15)

Proof. To prove this lemma, we only use

$$Q^4 = Q - Q''$$
 and $(Q')^2 = Q^2 - \frac{2}{5}Q^5$.

Proof of (12). By integration by parts,

$$\int Q^{r+3} = \int Q^{r-1}Q^4 = \int Q^{r-1}(Q - Q'') = \int Q^r + (r-1)\int (Q')^2 Q^{r-2}$$
$$= r \int Q^r - \frac{2(r-1)}{5} \int Q^{r+3}.$$

Thus.

$$\frac{2r+3}{5}\int Q^{r+3}=r\int Q^r.$$

Proof of (13). By similar arguments,

$$\begin{split} \int Q^{r+3} \left(\int_0^x Q^2 \right)^2 &= \int Q^{r-1} Q^4 \left(\int_0^x Q^2 \right)^2 = \int Q^{r-1} (Q - Q'') \left(\int_0^x Q^2 \right)^2 \\ &= \int Q^r \left(\int_0^x Q^2 \right)^2 + (r-1) \int (Q')^2 Q^{r-2} \left(\int_0^x Q^2 \right)^2 \\ &+ 2 \int Q^{r+1} Q' \int_0^x Q^2 \\ &= r \int Q^r \left(\int_0^x Q^2 \right)^2 - \frac{2(r-1)}{5} \int Q^{r+3} \left(\int_0^x Q^2 \right)^2 - \frac{2}{r+2} \int Q^{r+4}. \end{split}$$

Thus,

$$\frac{2r+3}{5} \int Q^{r+3} \left(\int_0^x Q^2 \right)^2 = r \int Q^r \left(\int_0^x Q^2 \right)^2 - \frac{2}{r+2} \int Q^{r+4}.$$

Proof of (14).

$$\begin{split} \int xQ^{r+3} \int_0^x Q^2 &= \int xQ^{r-1}Q^4 \int_0^x Q^2 = \int xQ^{r-1}(Q-Q'') \int_0^x Q^2 \\ &= \int xQ^r \int_0^x Q^2 + (r-1) \int x(Q')^2 Q^{r-2} \int_0^x Q^2 + \int Q^{r-1}Q' \int_0^x Q^2 \\ &+ \int xQ^{r+1}Q' \\ &= r \int xQ^r \int_0^x Q^2 - \frac{2(r-1)}{5} \int xQ^{r+3} \int_0^x Q^2 - \left(\frac{1}{r} + \frac{1}{r+2}\right) \int Q^{r+2}. \end{split}$$

Thus,

$$\frac{2r+3}{5} \int xQ^{r+3} \int_0^x Q^2 = r \int xQ^r \int_0^x Q^2 - \frac{2(r+1)}{r(r+2)} \int Q^{r+2}.$$

Proof of (15).

$$\int x^2 Q^{r+3} = \int x^2 Q^{r-1} (Q - Q'') = \int x^2 Q^r + (r-1) \int x^2 (Q')^2 Q^{r-2}$$
$$+ 2 \int x Q^{r-1} Q'$$
$$= r \int x^2 Q^r - \frac{2(r-1)}{5} \int x^2 Q^{r+3} - \frac{2}{r} \int Q^r.$$

Thus,

$$\frac{2r+3}{5} \int x^2 Q^{r+3} = r \int x^2 Q^r - \frac{2}{r} \int Q^r.$$

Proof of (9). This is [10, Claim 2.3]. We repeat it here for convenience.

We compute $\int Z_0 Q$, where $Z_0 = 3Q'' + 3V_0'' + 4Q^3V_0$. By $Q'' = Q - Q^4$, we get:

$$\int Z_0 Q = \int (3Q'' + 3V_0'' + 4Q^3 V_0) Q = 3 \int Q^2 - 3 \int Q^5 + \int V_0 (3Q'' + 4Q^4)$$
$$= 3 \int Q^2 - 3 \int Q^5 + \int V_0 (3Q + Q^4).$$

We compute the last term, integrating by parts:

$$\int V_0(3Q+Q^4) = -\int \left(\frac{1}{3}Q + \frac{3}{2}xQ'\right)(3Q+Q^4) = \frac{5}{4}\int Q^2 - \frac{1}{30}\int Q^5.$$

Finally, using Lemma A.1,

$$\int Z_0 Q = \frac{17}{4} \int Q^2 - \frac{91}{30} \int Q^5 = -\frac{1}{12} \int Q^2.$$

References

- J. L. Bona, W. G. Pritchard, and L. R. Scott, Solitary-wave interaction, Phys. Fluids 23 (1980), 438–441.
- [2] W. Craig, P. Guyenne, J. Hammack, D. Henderson, and C. Sulem, Solitary water wave interactions, Phys. Fluids 18 (2006), no. 5, 057106, 25.
- [3] E. Fermi, J. Pasta, and S. Ulam, Studies of nonlinear problems, I, Los Alamos Report LA1940 (1955); reprinted in A. C. Newell (ed.), Nonlinear wave motion, Lectures in Applied Mathematics, vol. 15, American Mathematical Society, Providence, R.I., 1974.
- [4] R. Hirota, Exact solution of the Korteweg-de Vries equation for multiple collisions of solitons, Phys. Rev. Lett. 27 (1971), 1192–1194.
- [5] P. D. Lax, Integrals of nonlinear equations of evolution and solitary waves, Comm. Pure Appl. Math. 21 (1968), 467–490.
- [6] Y. Li and D. H. Sattinger, Soliton collisions in the ion acoustic plasma equations, J. Math. Fluid Mech. 1 (1999), no. 1, 117–130.

- [7] Y. Martel, Asymptotic N-soliton-like solutions of the subcritical and critical generalized Korteweg-de Vries equations, Amer. J. Math. 127 (2005), no. 5, 1103-1140.
- [8] Y. Martel, F. Merle, and T.-P. Tsai, Stability and asymptotic stability in the energy space of the sum of N solitons for subcritical gKdV equations, Comm. Math. Phys. 231 (2002), no. 2, 347–373
- [9] Y. Martel and F. Merle, Refined asymptotics around solitons for gKdV equations, Discrete Contin. Dyn. Syst. 20 (2008), no. 2, 177–218.
- [10] _____, Description of two soliton collision for the quartic KdV equation, preprint, available at arXiv:0709.2672v1.
- [11] R. M. Miura, The Korteweg-de Vries equation: A survey of results, SIAM Rev. 18 (1976), no. 3,
- [12] M. Wadati and M. Toda, The exact N-soliton solution of the Korteweg-de Vries equation, J. Phys. Soc. Japan 32 (1972), 1403–1411.
- [13] N. J. Zabusky, Solitons and energy transport in nonlinear lattices, Comput. Phys. Comm. 5 (1973), 1-10.
- [14] N. J. Zabusky and M. D. Kruskal, Interaction of "solitons" in a collisionless plasma and recurrence of initial states, Phys. Rev. Lett. 15 (1965), 240–243.